

Statistical Mechanics

Thermodynamics and an Introduction to Thermostatistics, 2nd ed. by Herbert B. Callen

11/06

Outline:

- Goals of statistical mechanics
 - Formalism \Rightarrow Analytical engine
 - Examples
 - force-extension for DNA
 - DNA binding proteins
 - + gene regulation
 - Hemoglobin & Oxygen

Biophysics Lecture Series Talk

Jan. 23 - Steve Block

Macrostates

Microstates

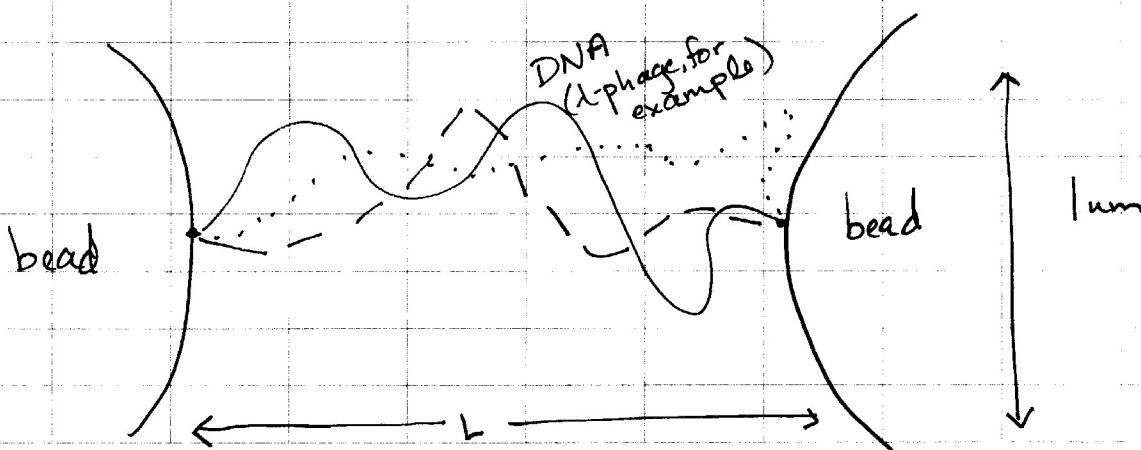
(example : all coordinates and momenta of all the microscopic particles)

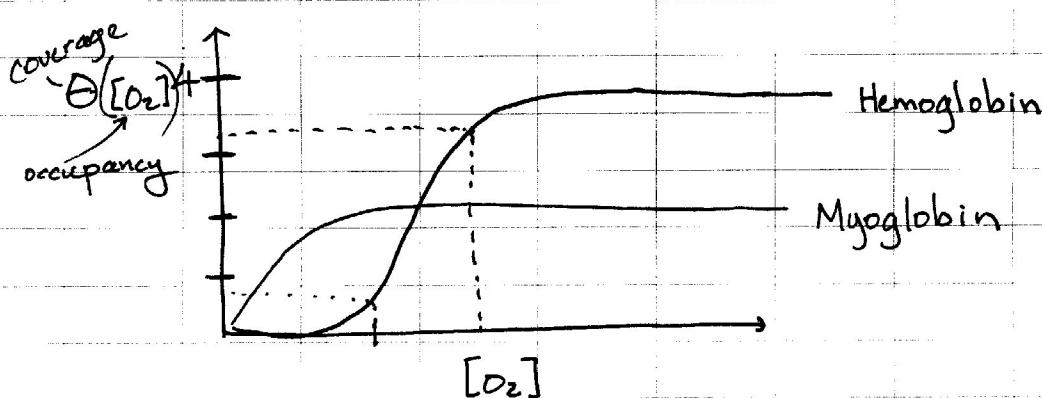
$$\{P_i, q_i\} = (P_1, P_2, \dots, P_N; q_1, q_2, \dots, q_N)$$

momentum
of molecule +

momentum
of molecule N

position of
molecules





A "Goal" of Statistical Mechanics

- find the probability of the microstates
 $\Rightarrow \text{prob(microstates)}$
e.g. gas : $p(\{p_i, q_i\})$

More formally :

let's label our microstates by their energies, E_n

$p(E_n)$ = probability of microstate with energy, E_n

Compute averages :

$$\text{e.g. } \langle E \rangle = \sum_n E_n p(E_n)$$

→ Example of a polymer

$$\langle \text{Length} \rangle = \sum_{\text{configurations}} p(\text{config}, f) \xleftarrow{\text{force}} l(\text{config})$$

Notion of Entropy

- Fundamental idea of thermodynamics & statistical mechanics

$$S(E) = k_B \ln \Omega(E)$$

↑ ↑
macroscopic microscopic

(Boltzmann's Law)

$\Omega(E) \equiv$ # of microstates
of energy E

2nd Law of Thermodynamics:

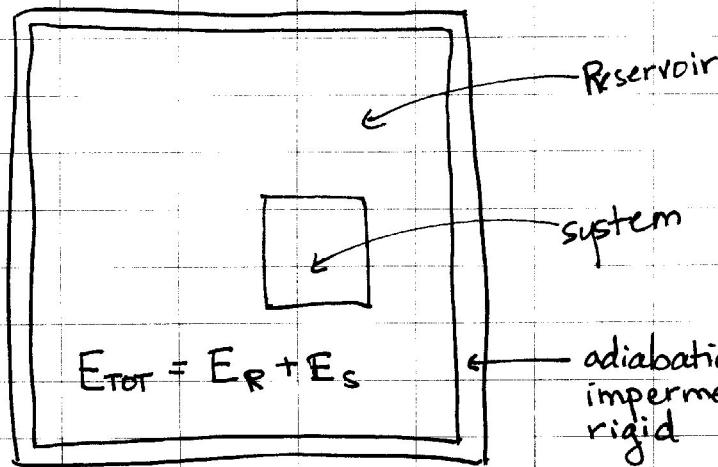
A closed system will reach the terminal privileged state of maximum entropy.

Josiah Willards Gibbs - translated the 2nd law of thermo. into mathematical calculations using DS.

State of maximum entropy is when $dS_{\substack{\uparrow \\ \text{equilibrium}}} = 0$ (Energy of system is constant)

* Boltzmann Distribution: $p(E_n) = \frac{e^{-E_n/kT}}{Z}$

$Z = \sum_n e^{-E_n/kT}$ = partition function
("analytical engine of statistical mechanics")



E_R = energy of reservoir
 E_S = energy of system

adibatic
impermeable
rigid] = the universe is isolated

$$S_{TOT}(E_{TOT}) = S_R(E_R) + S_S(E_S)$$

$$P(E_j) = \frac{\Omega_{RES}(E_{TOT} - E_j)}{\Omega_{TOT}(E_{TOT})} \quad \begin{matrix} \leftarrow \# \text{ of states where the reservoir has} \\ \text{energy of } E_{TOT} - E_j \end{matrix}$$

$\Omega_{TOT}(E_{TOT}) \leftarrow \text{total # of microstates available in the system + Reservoir}$

$$= \frac{e^{\ln \Omega_{RES}(E_{TOT} - E_j)}}{e^{\ln \Omega_{TOT}(E_{TOT})}}$$

$$= \frac{e^{S_{RES}(E_{TOT} - E_j)/k_B}}{e^{S_{TOT}(E_{TOT})/k_B}}$$

Aside #1 :

$$S_{RES}(E_{TOT} - E_j) = S_{RES}(E_{TOT} - \langle E \rangle + \langle E \rangle - E_j)$$

average energy of system

Aside #2 :
 $\frac{d}{dt} = \frac{ds}{dt}$

Taylor expansion
 $\rightarrow S_{RES}(E_{TOT} - \langle E \rangle) + \frac{ds}{dE}(\langle E \rangle - E_j)$

small parameter

$$P(E_j) = \frac{e^{S_{RES}(E_{TOT} - \langle E \rangle)/k_B}}{e^{S_{RES}(E_{TOT} - \langle E \rangle)/k_B}} e^{\frac{1}{k_B T}(\langle E \rangle - E_j)} \frac{e^{S_{SUS}(\langle E \rangle)/k_B}}{e^{S_{SUS}(\langle E \rangle)/k_B}}$$

$$S_{TOT}(E_{TOT}) = S_{RES}(E_{TOT} - \langle E \rangle) + S_{SUS}(\langle E \rangle)$$

$$\Rightarrow p(E_j) = \text{const. } e^{-E_j/k_B T} = e^{-E_j/k_B T} e^{\frac{1}{k_B T} (\langle E \rangle - S_{\text{sys}}(\langle E \rangle) T)}$$
$$= e^{-E_j/k_B T} e^{F/k_B T}$$

Normalization : $\sum_n p(E_n) = 1$

$$\sum_j e^{\beta F} e^{-\beta E_j} = 1 \Rightarrow e^{\beta F} \sum_j e^{-\beta E_j} = 1$$

$$e^{\beta F} = \frac{1}{\sum_j e^{-\beta E_j}} = \frac{1}{Z}$$

$$\Rightarrow Z = \sum_j e^{-\beta E_j} \quad \leftarrow \beta = \frac{1}{k_B T}$$

Know these:

$$(1) \quad p(E_j) = \frac{1}{Z} e^{-\beta E_j} \quad \text{Boltzmann Distribution}$$

$$(2) \quad Z = \sum_j e^{-E_j/k_B T} \quad \text{Partition Function}$$