

Energetic gradients emerge in developing motor-microtubule structures

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This manuscript was compiled on May 18, 2026

Living matter produces a variety of beautiful spatiotemporal structures and patterns that are not enduringly present in their nonliving counterparts. These ordered, non-equilibrium steady states are often sustained through the consumption of energy. Here, we investigate the energetic cost of assembling an ordered aster from an initially disordered, uniform mixture of cytoskeletal microtubules and kinesin motors. Using a calibrated fluorescent ATP reporter, we measure reproducible radial ATP gradients on scales of tens of microns that establish within, and persist over, tens of minutes, alongside coupled spatial gradients in motor density. These appreciable gradients are predicted by a reaction-diffusion model that acknowledges the localization of ATP consumption to regions where both molecular motors and microtubules are sufficiently abundant to encourage consumption, as confirmed by finite element modeling. With our results, we compare the power per volume required by our cytoskeletal networks with the known power per volume expenditure in cells. Comparison of our measured results with estimates of the dissipative processes available to motor-microtubule mixtures leads to the hypothesis that maintaining spatial motor gradients dominates the energetic demand in this system. Our direct quantification of energetic fluxes across space unlocks future explorations of what steady states are accessible to cells, and how the cytoskeleton drives broad spatial organization.

nonequilibrium|energetics|active matter|cytoskeletal organization|motor-microtubule assemblies

A key driver of the rich patterns in living matter is a steady investment of energy. For example, the establishment of morphogen gradients in developmental patterning arises from its synthesis, spatial redistribution, and degradation (1–4). Similarly, cytoskeletal motor systems hydrolyze adenosine triphosphate (ATP) and other triphosphates to achieve processes ranging from intracellular transport, to cell motility, to chromosome segregation. Motivated by these processes, we were inspired to develop a physical understanding of the connection between energy fluxes and the emergence of biological order in space and time in the context of the particular example of microtubule-motor assemblies.

The energetic basis of the processes within living cells are based upon a few fundamental energy currencies, which can be thought of as biological batteries. This metaphor is useful because it reminds us that batteries are indifferent to the particulars of what they are wired up to—they can drive anything from the light in a flashlight to motorized toys. Biological processes are powered by several key biological batteries including membrane potentials, redox reactions and trinucleotide hydrolysis. Indeed, for the molecular motor driven reactions that power the structures of the cytoskeleton, ATP and GTP hydrolysis are central. Thus, we were curious about how ATP consumption in space and time drives the dynamics of structure formation.

Significance

How much energy do organisms pay to form and maintain their organizing biochemical patterns? Existing measurements of cellular metabolism and energy expenditures largely resolve net or supply-side biochemical fluxes, without spatial information, impeding the study of this basic question. Here, we develop an experimental approach to directly measure the distributions of biochemical energy that respond to power expenditures of cytoskeletal motor-microtubule networks as they form aster structures, reminiscent of those found in the mitotic spindle. As these structures self-assemble, calibrated readouts in real molecular units register large, reproducible, and long-lived gradients of ATP. We interpret these measurements by developing theory to account for the functional destinies of energy expenditure. These advances clarify outstanding questions of energy in living matter.

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A.I.D., G.L.S., H.J.L., R.B., M.T., and R.P. designed research; A.I.D., G.L.S., H.J.L., and B.N., performed research; A.I.D., G.L.S., H.J.L., B.N., M.A., S.H., H.W.Ch.P., and R.B. contributed new reagents/analytic tools; A.I.D., G.L.S. and M.A. analyzed data; and A.I.D., G.L.S., and R.P. wrote the paper.

The authors declare no competing interest.

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123 Recently, it has become possible to measure the total
124 energy consumption of both living organisms and the
125 molecular machinery that drives them (5–9), joining
126 theoretical methods to infer energy expenditures from
127 observables (10). Despite the foundational value of these
128 investigations, to date almost all these measurements
129 report such metabolic expenditures as spatial and tem-
130 poral averages. These empirical limitations mean that
131 putatively profound heterogeneities in expenditures across
132 regions and processes of the cell remain unknown and
133 unmeasured.

134 How significantly does dissipation localize in space
135 and time among biomolecular assemblies? How big and
136 sustainable are bulk energetic and dissipative gradients
137 sculpting living matter? To address these open questions
138 and complement earlier foundational studies, we set
139 out to make precise, real time measurements of energy
140 abundance and consumption resolved on the micron
141 scale. In particular, here we report the visualization
142 of spatial ATP concentration gradients across cytoskeletal
143 networks, giving insight into how structure, composition
144 and morphology drive energy dissipation.

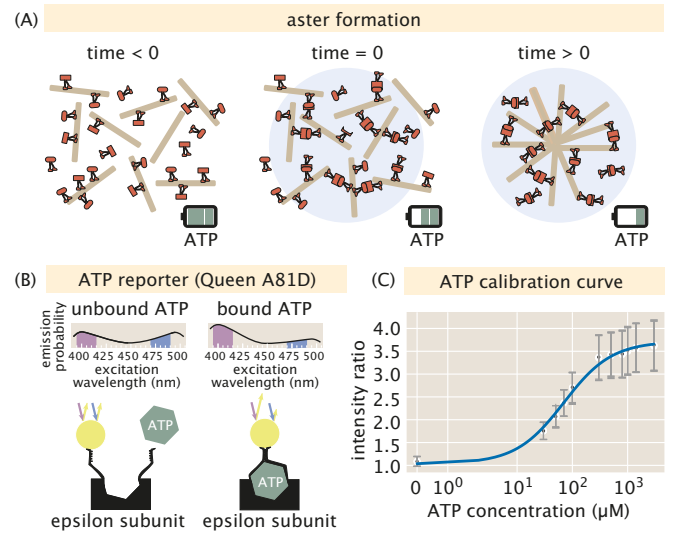
145 We use motor-microtubule assemblies as a highly
146 controllable and tunable system due to their minimal
147 components and self-organizing properties *in vitro* as
148 well as their biological ubiquity (11–16). An abundance
149 of work has established that connected dimeric motor
150 domains can cross-link microtubules, creating ordered
151 networks (11, 12, 17, 18). To control the position, size
152 and start of microtubule cross-linking, we optogenetically
153 link motor proteins together as shown in Figure 1, as
154 previously developed in our labs (18). Motors harness
155 energy to drag microtubules into ordered structures
156 by hydrolyzing an ATP molecule for each step they
157 take along the microtubule. We measure the energy
158 consumed by the motors throughout space and time
159 using a fluorescent, QUEEN-based ATP probe (19). We
160 illustrate the ratiometric probe mechanism shown in
161 Figure 1(B).

162 In addition, to complement the measurements and
163 to provide a framework for understanding them, we
164 examine a reaction-diffusion model that describes the
165 emergent ATP spatiotemporal gradients and explore the
166 implications of that model with finite element calculations.
167 Combining our experimental and theoretical results, we
168 determine that the measured network formation power
169 is indeed many orders of magnitude greater than the
170 theoretical power of equilibrium processes.

171 Results

172 Direct measurement and visualization of emergent ATP gradients.

173 The key elements of our experimental design are shown in
174 Figure 1. As noted above, using spatially and temporally
175 controlled illumination, we can generate patterns such
176 as the radially symmetric aster shown in the schematic.
177 Our principal experimental goal is to measure the rate of
178 consumption of ATP as a function of position and time, a
179 goal that is realized by using the fluorescent, ratiometric
180 ATP reporter (19) depicted in Figure 1(B). The probe
181 mechanism creates a change in the protonation state (20)
182 of the fluorophore when ATP binds (21), triggering a shift
183 in the fluorophore’s absorption spectrum (22). By using
184 known standards, as shown in Figure 1(C), we construct
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Fig. 1. Schematic of the experimental system used to measure spatiotemporal evolution of ATP. (A) The formation of an aster using light activated motor dimerization. Before light activation, motors independently walk on microtubules, uniformly hydrolyzing ATP. At $t = 0$ a circular light pattern is projected onto the sample. Motor proteins inside the illuminated region dimerize, crosslinking microtubules. As time elapses, microtubules are dragged into an aster resulting in spatially heterogeneous depletion of ATP. (B) The binding mechanism for ATP to the ATP probe. Binding ATP to the probe causes the number of emission counts due to an excitation of 405 nm light to increase, while emission counts at 480 nm light excitation decreases. By comparing the ratio of the emission counts at 405 nm and 480 nm light excitations, the concentration of ATP can be inferred. (C) A calibration curve mapping known ATP concentrations to fluorescent light intensity ratios. Each black circle represents the mean ratio value for a given image and gray error bars report the standard deviation of the image.

220 a calibration curve that permits us to measure the ATP
221 concentration in a given spatial region. Given that the
222 characteristic scale of ATP concentrations in our exper-
223 iments are of order a few hundred μM , our calibration
224 curve, in Figure 1(C), shows that the ATP reporter is
225 sensitive in precisely the concentration range needed to
226 spatially resolve ATP consumption in the system. The
227 Supplemental Information, provides a detailed description
228 of how we handled uneven illumination (section S2.B) and
229 photobleaching (section S2.C), a prerequisite to generate
230 such a calibration curve. There we also describe how we
231 used two-dimensional images to infer the properties of our
232 three dimensional system (section S7).

233 This measurement scheme equips us to simultaneously
234 resolve the quantity of both ATP (the “fuel”) and molec-
235 ular motors (the “consumer”) over space and time while
236 asters form, as shown for representative time courses in
237 Figure 2 (Top) and (Bottom). Local motor concentrations
238 in space are inferred from fluorescence footage using a
239 separate calibration curve, as discussed further in the
240 Supplemental Information (section S2.D). As motors
241 step along and exert torques on microtubules, they
242 accumulate in the centers of assembling asters, reaching
243 a peak concentration of $\approx 3 \mu\text{M}$, creating self-organized
244 polar order and material flow, as witnessed in the time
245 course of Figure 2(Top). Concurrently, our measurements
246 report how an initially uniform concentration of ATP
247 is steadily reduced in time, depleting from an initial
248 concentration of $500 \mu\text{M}$ to around $350 \mu\text{M}$ in the first
249 four minutes, to $\approx 100 \mu\text{M}$ in about twenty minutes, as
250 shown by Figure 2(Bottom). Importantly, the depletion

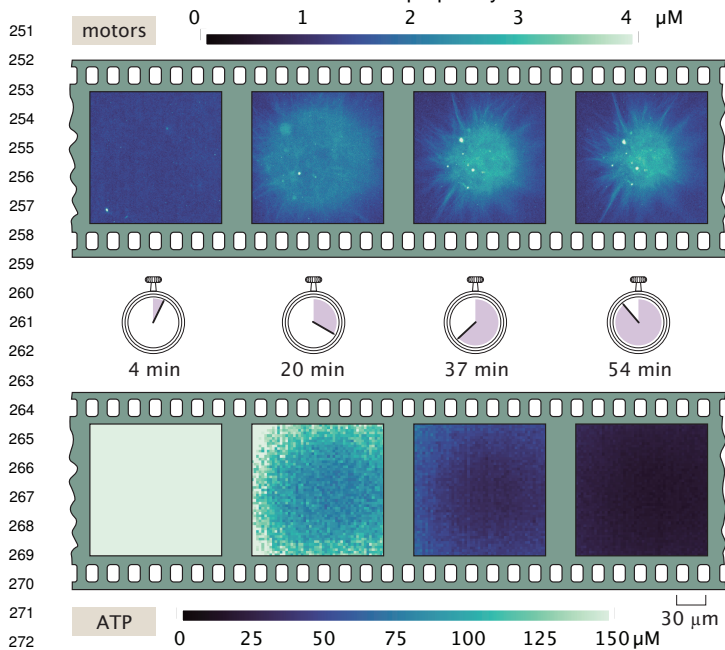


Fig. 2. Experimental measurements resolve coupled gradients of motors and ATP across space and time. (Top) Experimentally measured spatial distributions of molecular motors and (Bottom) ATP over four time points during the self-organization of an aster. As time evolves, motor proteins concentrate near the aster center; a coupled ATP gradient develops, with greatest depletion in the aster's center where motors are most abundant.

of ATP is manifestly nonuniform over space, forming a steepening gradient of ATP, with less ATP in the aster center than at its periphery. These concentration fields are displayed as radial profiles in Figure 3(A) and Figure 3(B), registering rich time and space dependencies. Specifically, conservatively summarizing ATP gradients by their change in concentration across a one hundred micron radial region near the edge of an aster (between radial positions $r = 200 \mu\text{m}$ near the edge and $r = 100 \mu\text{m}$) registers gradients of $0.5 \mu\text{M}/\mu\text{m}$ to $1 \mu\text{M}/\mu\text{m}$ across replicate asters (see Supplementary Material section S3).

In general, the uncertainties in such interesting physical scales (such as gradients and dissipations) reported in this paper are most earnestly captured by their aggregate variations across many separate aster formation experiments, as detailed in the Supplementary Information, section S3, Figures S25-S28.

Appreciating the biological magnitude of gradients. Are the spatial gradients in ATP reported by our calibrated measurements in Figure 2 and Figure 3 with a characteristic scale of $\approx \mu\text{M}/\mu\text{m}$ biologically meaningful? To develop intuition for that question, we examine how the measured ATP concentrations vary relative to the K_M of the motors. At intermediate times, as highlighted by Figure 4(A), ATP drops from a high value of $\approx 100 \mu\text{M}$ ATP that is several times larger than the characteristic concentration scales governing motor stepping activity (namely the enzymatic Michaelis-Menten constant $K_M \approx 23 \mu\text{M}$ ATP, as reported by steady-state ATPase assays (23, 24)). Further, at sufficiently late times, ATP depletes sufficiently so as to be above this characteristic K_M at the edge of the aster but below this K_M in the center of the aster (as illustrated by the 37 min pink-fuchsia inset of Figure 3(B)).

Thus, the spatial variation of ATP inside a single aster lies squarely in the functional range that controls whether or not motors have enough ATP to move close to their maximum stepping rates.

Another way to size up these gradients is to compare them to other important gradients as we detail in the Supplementary Information section S10. Figures 4(B), (C), and (D) provide several candidate reference points. In Figure 4(B), we illustrate a classic experimental study (25) on bacterial chemotaxis in which bacteria navigate gradients of approximately $30 (\text{molecules}/\mu\text{m}^3)/\mu\text{m}$ (25). As seen in Figure 4(C), the protein kinase Pom1 in fission yeast is part of the signaling pathway that controls entry to mitosis, Pom1 localizes in the poles of the yeast cell (26). We estimate the gradient in these yeast cells is approximately $3 (\text{proteins}/\mu\text{m}^3)/\mu\text{m}$. As another example, we consider the transcription factor *Bicoid*, which famously contributes to delivering the positional information that organizes the fruit fly *Drosophila*'s body plan, as shown in Figure 4(D) (2, 27–29). This molecular gradient develops with high precision across embryos, and robustly adopts a characteristic profile with a gradient of $\approx 0.2 (\text{molecules}/\mu\text{m}^3)/\mu\text{m}$. Thus we see that the maximal ATP gradients we measure are ten to a few thousand times steeper than the chemoattractant, Pom1 and *Bicoid* gradients used to drive bacterial chemotaxis, cellular division and sculpt embryos, respectively. (Further numerical discussion, including complementary notions of steepness, is found in the Supplementary Information section S10.) Accordingly, these measurements attest that such gradients could have biologically salient impacts.

Mapping power consumption in space and time. Another way of visualizing the results of our measurements is to take the ATP data from successive instants and convert it to a power. As shown in Figure 5, the power can be evaluated directly in units of ATP/s, revealing power of order $\text{few} \times 10^8 \text{ ATP/s}$. This magnitude of power expenditures is reproducibly observed across replicate asters as described in detail in the Supplementary Information, section 3B.2, Figures S25 and S29; peak dissipation rates usually occur at early times, and vary between $\approx 3 \times 10^8 \text{ ATP/s}$ and $\approx 8 \times 10^8 \text{ ATP/s}$ across separate self-organizations of distinct asters. We interpret this few-fold variation in the dissipations across replicate aster experiences as representative of aggregate uncertainty in these scales (see SI).

Note that these changes in the ATP present in an aster over time technically reflect both a consumption power inside the aster's volume, and any (not explicitly measured) replenishment flux of ATP into the aster due to diffusion or fluid flow. In the present setting, such replenishment fluxes should counter consumption: since ATP increases with radial position (including at the aster's boundary), average diffusive flux directs inwards (by Fick's law), and net fluid flow is also inwards in this contractile setting (see Supplementary Information section S2J). Accordingly, the net ATP consumption rates reported by measurements here may be regarded as reasonably close lower bounds to full underlying consumption rates.

These real thermodynamic units of our calibrated measurements allow us to compare our total measured power with estimates of what the energy from hydrolysis events is used to pay for. In addition, we can directly

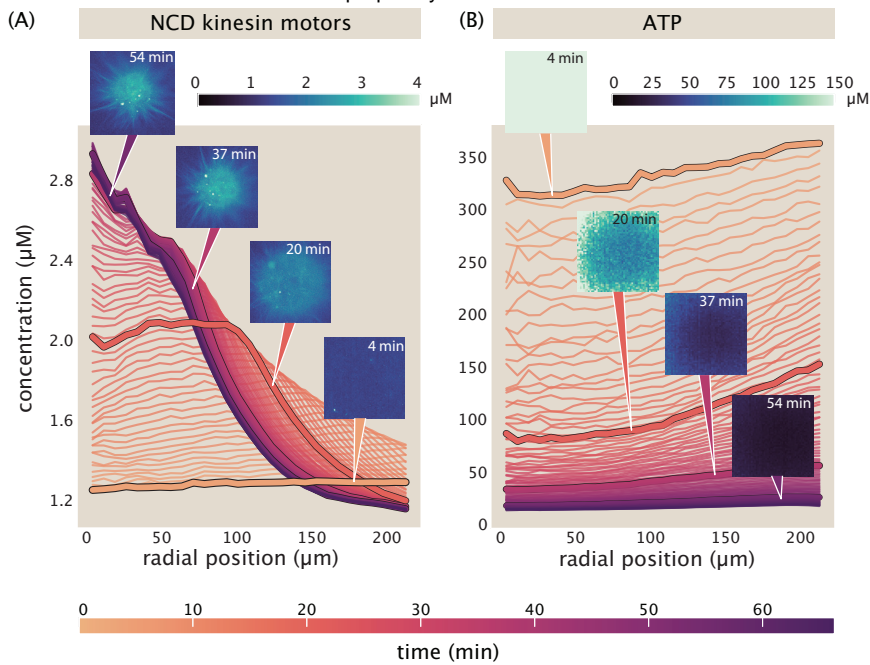


Fig. 3. Measurements of developing ATP and motor gradients with respect to radial position in asters. (A) Azimuthally-averaging the data shown in Figure 2(Top) gives radial concentration profiles of motors. (B) Similarly, azimuthal averaging of the data shown in Figure 2(Bottom) yields the radial concentration profiles of ATP. These gradients reveal clear, rich nonuniformities over time and space. Further replicate aster experiments show largely coherent behaviors, with scales of variation across replicate experiments visible in the Supplementary Information section S3B.1

contrast these measurements with power expenditures reported in far different biological contexts. We return to reckon with what sets the scale of these measured power values in the final *Results* section. Next, we ask how one might develop quantitative intuition for the distribution of the ATP in the aster in both space and time.

Understanding measured gradients with reaction-diffusion modeling.

The profiles revealed in Figure 6 characterize the radial and temporal dependence of both the motors and the ATP. We present a dynamical equation in Figure 6(A) which acts as a minimal model to describe the resultant ATP profiles through space and time. The rate of change of ATP in a small material volume element can be attributed both to ATP molecules entering and leaving that small region (diffusion) and to the hydrolysis of those ATP molecules by molecular motors that are in the material volume element of interest (reaction). The reaction term is defined as the product of the hydrolysis rate of a motor protein, the motor profile (which gives the density of motors in that material volume element), the bound fraction of motors to microtubules, and the probability an ATP is bound to a motor. The bound fraction of motors to microtubules is described by a kinetic model of binding as derived in the Supplementary Information section S8A. The probability an ATP is bound to a motor protein is determined by Michaelis-Menten-like dynamics for the rate of ATP consumption by motors. However, the denominator of that term also includes terms that reflect competitive inhibition of the reaction due to ADP and P_i . Note that there is another more pernicious dynamic taking place during our experiments, namely, photobleaching. Section S6 in the Supplemental Information describes how we measure and account for photobleaching, including through finite element simulations of a diffusion-photobleaching equation in section S6F.

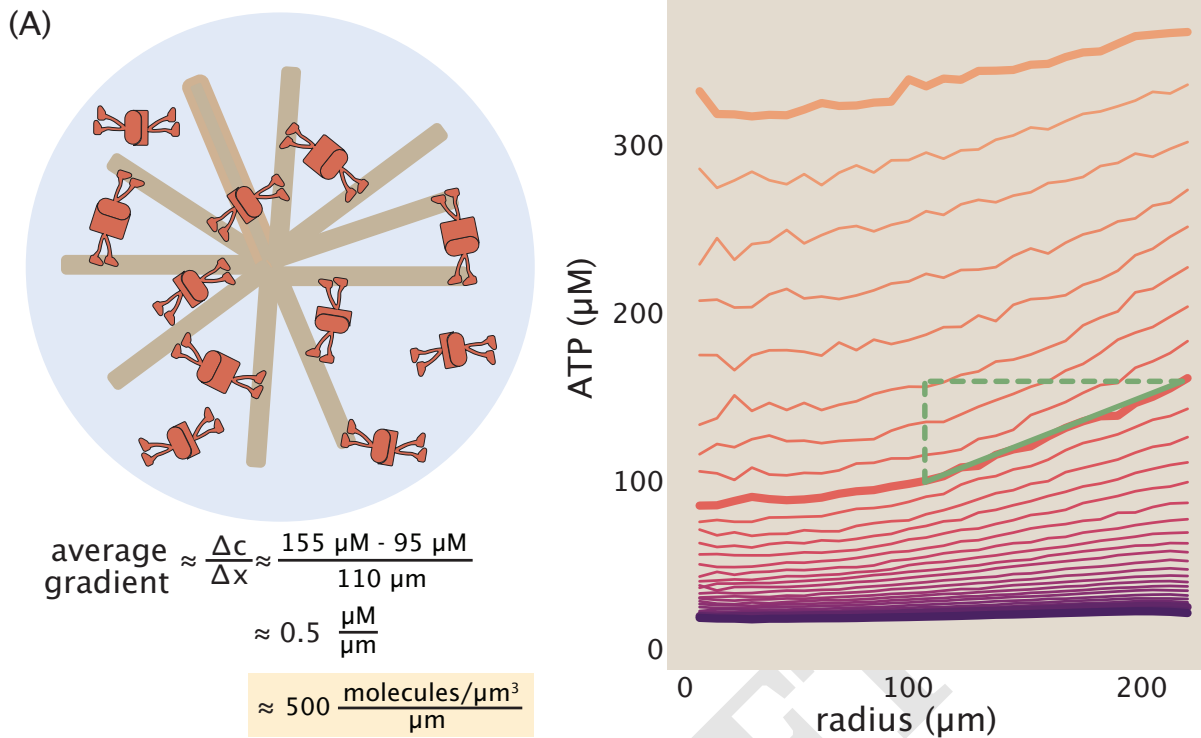
One powerful way to analyze the solutions to equations such as that shown in Figure 6(A) in diverse geometries is by appealing to numerical methods. In our case, we used the finite element method to compute the space-

time history of ATP as shown in Figure 6(C). To generate these plots, we input trial functions for the time-dependent motor and tubulin profiles, which are based on measured data. The form of these profiles are discussed in detail in the Supplementary Information section S4A.4 and are depicted in Figure 6(B), along with the resultant bound fraction of motors and bound motor concentration across time and space. To assess the quality of our simulation, we consider two semi-quantitative checks: does the global concentration of ATP deplete on a similar time scale to that measured in the experiment and is the size of the maximum gradient comparable to the maximum gradient in the experimental profiles? In Figure 6(C), the left plot illustrates these two dynamical metrics. The second graph to the left plots the experimental ATP profiles, which we note show a maximal slope of $\approx 0.6 \mu\text{M}/\mu\text{m}$ and globally depletes in ≈ 66 minutes. The two plots on the right show the results of our simulation. The second plot from the right uses our best estimates for the parameters that we believe apply to our experiment, which we expound on in the Supplementary Information section S4A. For this particular choice of parameters, the simulated profiles have a much shallower gradient and globally deplete ATP slower. This inspired us to systematically sweep the parameter values to find those that more faithfully reproduce the measured profiles. We find in the right most plot that lowering the diffusion constant by half, raising the motor hydrolysis rate by 50% and increasing the maximum tubulin concentration by a factor of 3 results in an ATP depletion profile that better matches our data. Here we see a maximum gradient of $0.5 \mu\text{M}/\mu\text{m}$ and global ATP depletion in 58 minutes. Note that these alternative parameter values are still within the uncertainty of knowledge of their values.

Note that our goal with these simulations is not to fit data curves, but rather to examine plausible behaviors for realistic parameter choices. With reasonable parameter choices, within a factor of a few of our assumed parameter values, we find gradients in ATP of the right scale appear in our simulation. While we attempted to characterize

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this study's ATP gradients



scales of biologically relevant gradients

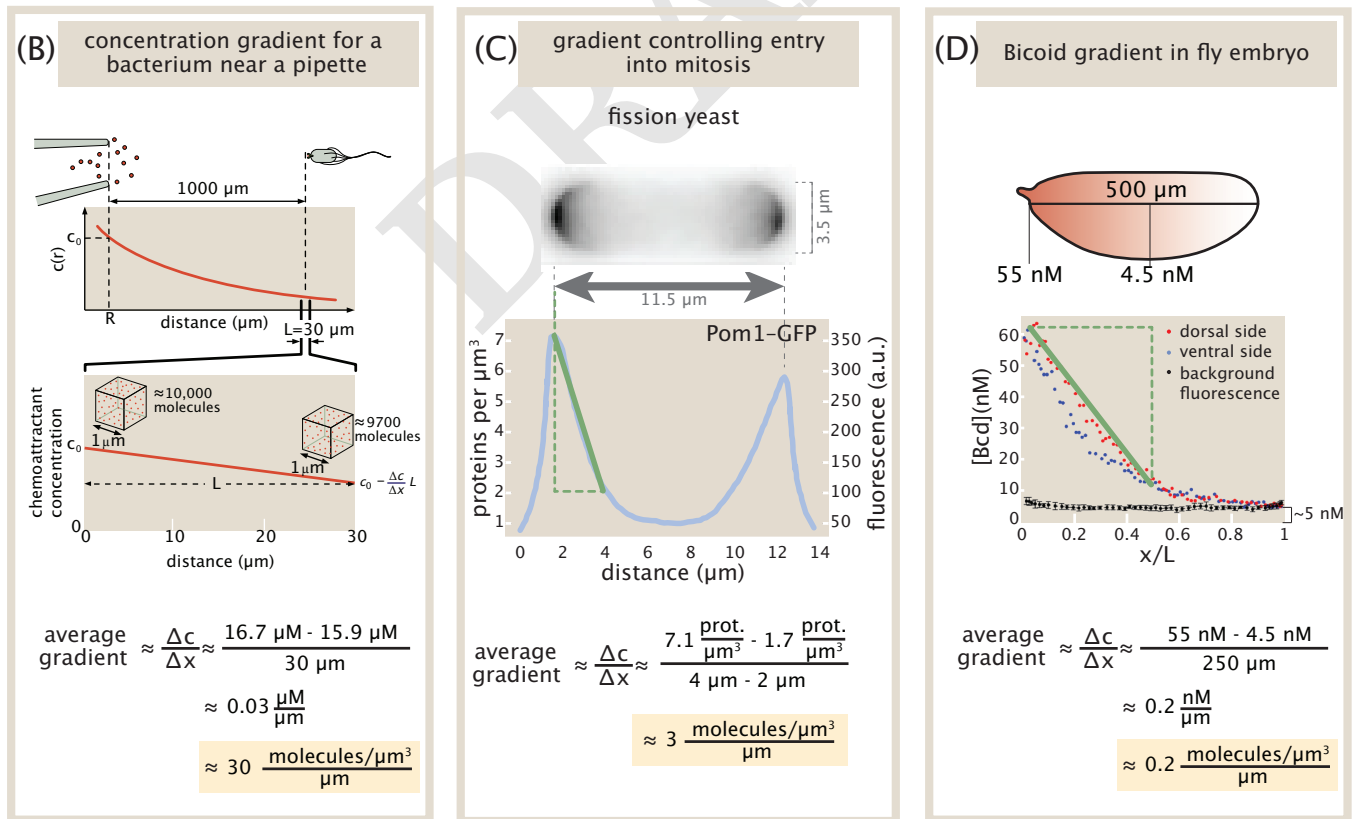


Fig. 4. Appreciating the relative scales of biologically developed gradients. (A) Our measured ATP gradients experience approximately a factor of two depletion across the scale of the aster length. (B) Bacterial cells respond to gradients of chemoattractants (Plot based on V. Sourjik et al. (25)). (C) In fission yeast, the protein kinase Pom1 localizes at the cell poles creating steep gradients on each side of the cell. (Plot adapted from J. Moseley et al. (26).) (D) The *Bicoid* gradient organizing early fly development manifests the majority of its decrease on the scale of half its body length. (Plot adapted from T. Gregor et al., (29).)

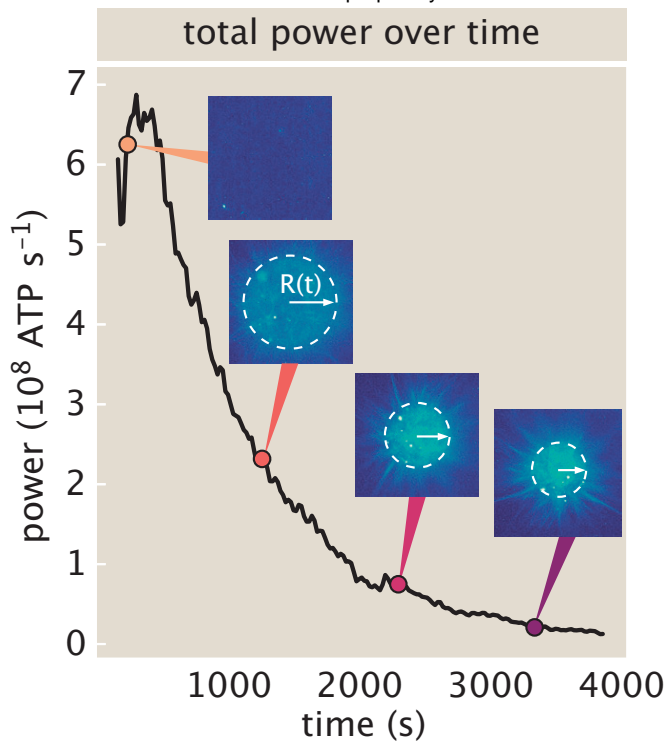


Fig. 5. Experimental readouts of total power during aster formation. Time derivatives of ATP reflect local and global power consumption, acknowledging the imaging geometry. The number of ATP molecules consumed per time over the whole aster volume significantly changes over time. Under these conditions ($1.2 \mu\text{M}$ motor proteins and $500 \mu\text{M}$ initial ATP), the magnitude of power tends to follow the size of the aster. Initially, while the aster is most rapidly contracting, ATP is consumed most rapidly, then approaches a baseline power level once the aster is no longer dramatically changing in size. This trend is visible in the majority of aster replicates which can be seen in the Supplementary Information Figure S29.

the most significant processes impacting the rate of ATP consumption, we leave the open questions of what other mechanisms may impact the shape of gradients. These may include fluid flow, motor cooperativity, motor jamming, motor detachment, and motor force-stalling, to name just a few.

More generally, the adequate accord of such a reaction-diffusion model for how gradients develop in ATP with measurements attests to general principles for how molecular gradients can form in space. This mechanism differs in fundamental conceptual ways from the simple and classical mechanism that develops the *Bicoid* gradient in *Drosophila*. There, in *Drosophila*, some reaction (degradation) rate k uniform in space conspires with a nonuniform boundary condition in space (specifying a uniform production in time at one end of the embryo), and competes with diffusion at coefficient D to set a steady-state gradient with a characteristic length scale $\lambda \sim \sqrt{D/k}$. However, applying identical logic to the present motor-microtubule system would be hasty (and would anticipate an incorrectly and negligibly-long analogous gradient length scale).

Instead, we contend that gradients in the present active matter system ensue very differently. Here, gradients in ATP develop when its hydrolysis reaction grows sufficiently nonuniform in space (thanks to the underlying self-organization of motors and microtubules), against uniform boundary conditions in space. Specifically, gradients

ensue from sufficiently dramatic localization of hydrolysis rates, whose effects accumulate for long enough times compared to the characteristic diffusive smoothing times that appreciable gradients sustain over pattern formation. Indeed, this connection between a field's spatial gradient and that accumulated over time in a consumption rate field can be expressed as a precise mathematical statement: the steepest gradient in concentration is at most the maximal gradient in hydrolysis rates accumulated over earlier times (see Supplementary Information, section S13). Such concentration gradients thus testify to reaction rate gradients. In sum, then, the precise competition of characteristic physics which govern whether a gradient appears, and its quantitative extent, differs from other celebrated classical gradients organizing biology. That an emergent patterning of reaction rates itself controls the appearance of sustained gradients may be a principle that operates beyond the specific motor-microtubule setting explored in this work.

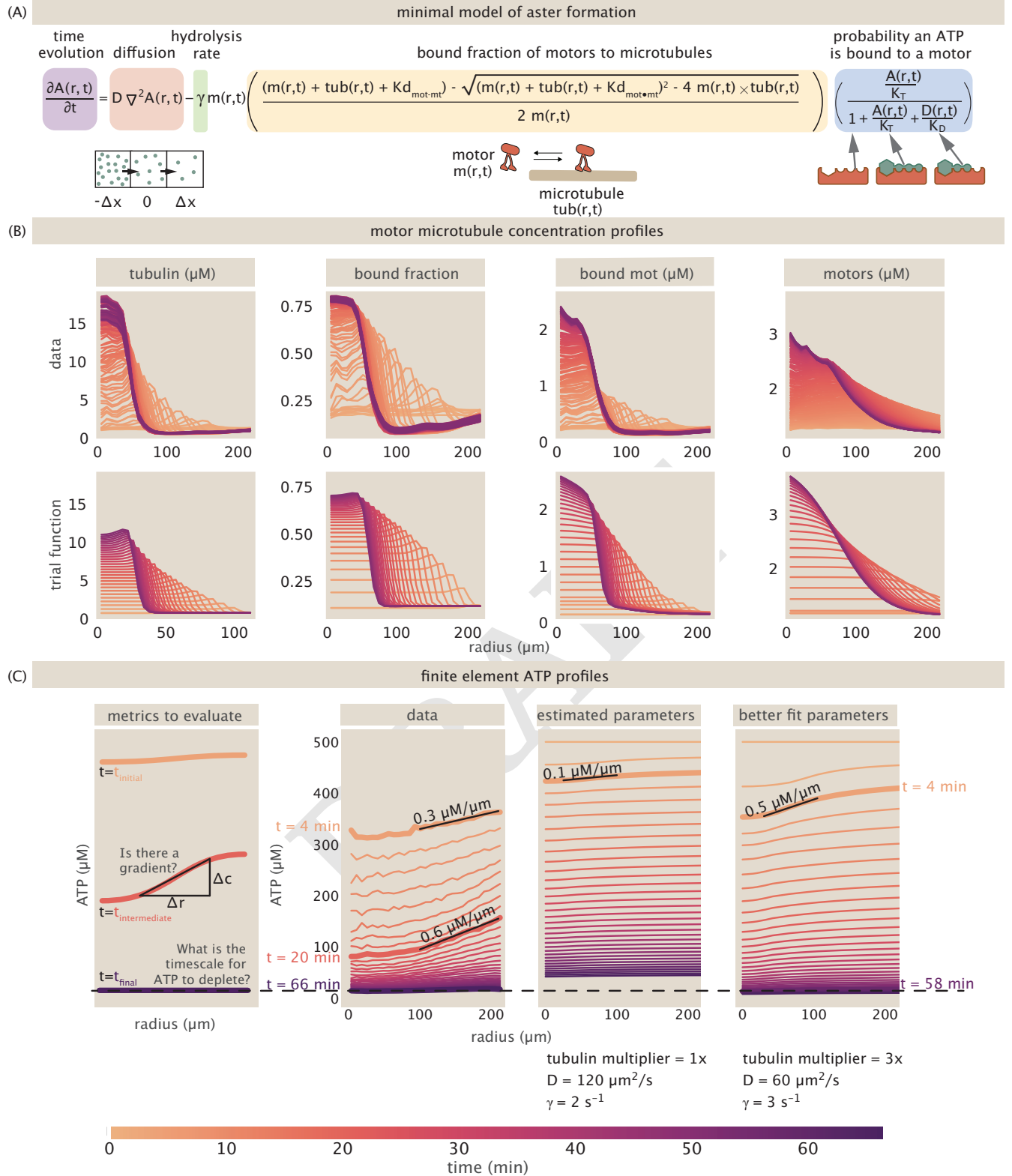
Comparing and appraising underlying physical origins for measured power. Provoked by our measurements and these analyses in the preceding sections, a fundamental and urgent question arises. What does that ATP hydrolysis “pay for?” Of course, mechanistically, we know that ATP is being consumed by motors as they carry out their walk along microtubules. But here we mean it differently and functionally. That energy is dissipated through elementary processes such as ordering and contraction of the microtubule network. How much energy do these processes cost?

Microscopically, a huge variety of dissipative processes are taking place during the microtubule-motor rearrangements attending aster formation. As shown in Figure 5 the power varies considerably at different stages of the aster formation process. To that end, we explore the power associated with a variety of processes that we imagine are taking place concurrently and would cost different amounts at different stages of the aster formation process.

Given the measured power, we were intrigued to compare it to the power associated with a variety of elementary dissipative processes that take place during aster formation as shown in Figure 7. For example, as is evident from the radius as a function of time in Figure 5, the volume of the microtubule aster is decreasing over time. As shown in Figure 7(A) and described in detail in the Supplementary Information section S14C, we can perform a simple estimate of the power associated with this contraction as the pressure-volume work done divided by the elapsed time. We find that the pressure-volume power is five orders of magnitude smaller than the measured power. As shown in Figure 7(B), a second dissipative process is the frictional sliding of the various microtubules during the contraction process. A naïve estimate is obtained by replacing a given microtubule by a corresponding sphere of the same dimensions and to work out the Stokes drag (see Supplementary Information section S14E). As in the case of the pressure-volume power, this results in a power that is five orders of magnitude smaller than the measured power of Figure 5. As discussed in the Supplementary Information section S14E.1, a better estimate can be made in which the microtubule is treated as a rod rather than a sphere and in this case the computed power is even smaller. It is possible that crowding

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Fig. 6. Minimal mathematical model and finite element simulations of ATP concentration in space and time. (A) The reaction-diffusion equation used to simulate ATP concentration is written with illustrations of the diffusive term, the binding states of motors to microtubules, and the binding states of ATP (A) and ADP (D) to the motor protein. (B) The distribution of microtubules and motors as a function of distance from the center of the aster as a function of time show the development of a gradient over time. The bottom row shows the empirical and approximate fits used as input into the finite element calculations. The key point of the motor and microtubule temporal distributions is that they are used as input into the consumption-diffusion model for the ATP. (C) We compare the gradients produced by finite element simulations on two metrics: how quickly the ATP depletes and the slope of the largest gradient. Next to the data, we plot the simulation using our best estimated parameters, namely a diffusion constant of 120 $\mu\text{m}^2/\text{s}$ and a hydrolysis rate of 2 s^{-1} . Additionally, the inputted tubulin profile is the trial function profile in panel (B). The right most plot shows parameters that better satisfy our evaluation metrics using a diffusion constant of 60 $\mu\text{m}^2/\text{s}$, a hydrolysis rate of 3 s^{-1} , and multiplying the maximum value of the tubulin profile by a factor of 3.

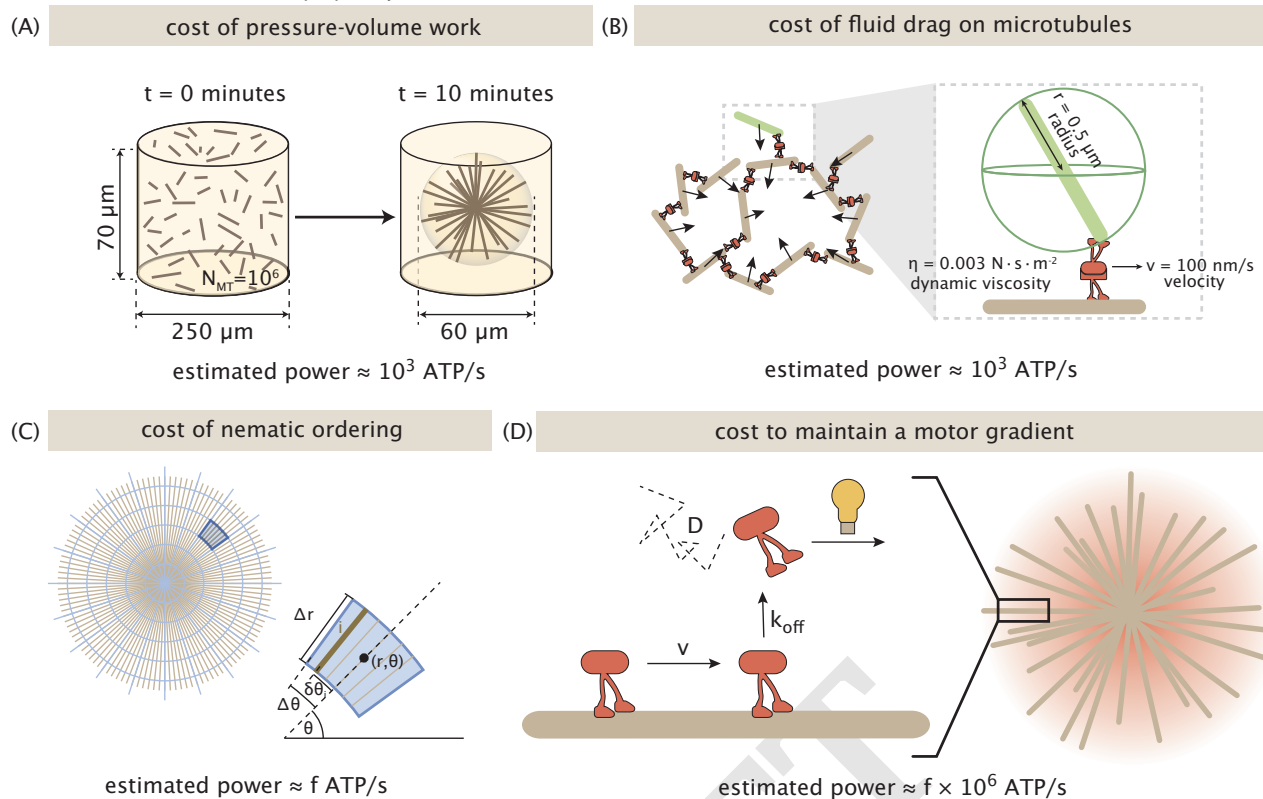


Fig. 7. Mechanistic processes taking place during aster formation and their estimated power. Each schematic considers a different dissipative processes that occurs during the formation of an aster. (A) Schematic of the power of compressing an "ideal gas" of microtubules. (B) The power of dragging microtubules through a viscous medium. (C) Power estimate for inducing nematic ordering in a random array of microtubules. (D) The power to maintain a concentration gradient of motors in the aster.

effects could amend these estimates. Another approach to estimating the power is offered by field theories of nematic ordering in which the state of the system is characterized by the spatially varying tensor $Q_{ij}(\mathbf{r}, t)$. This estimate is trickier to make since the parameters in such a field theory of motor-microtubule systems are not well known. Nevertheless, as seen in Figure 7(C) (and described in more detail in the Supplementary Information section S14F), the power we estimate associated with such ordering is many orders of magnitude smaller than the measured power. The final dissipative process highlighted in Figure 7(D) is that of building and maintaining a nonequilibrium gradient of motors radially outward from the center of the aster. One way to think about such a gradient is that if there were not some active transport carrying motors towards the aster center, then diffusion would smooth out that gradient. As we show in Supplementary Information section S14G, there is a well-defined prescription for estimating the power to maintain such a gradient using statistical physics; we find that this estimated expenditure at late times is roughly an order of magnitude lower than our measured powers at late times. This interesting result suggests the hypothesis that a significant fraction of the ATP hydrolysis consumed by the motors is "spent" to build and then maintain this gradient.

Discussion

It is practically a cliché to note that living organisms are "out of equilibrium." And yet, because of separation of time scales, equilibrium ideas are often useful in a variety of

nonequilibrium settings. To render our approach to these problems more precise, it is useful to measure the rate at which energy is being consumed to maintain systems in these nonequilibrium states. One useful way to think about such dissipative processes in living organisms is to recast them not as a cost, but rather as a mechanism.

Recent work has made exciting strides in characterizing energy expenditure in a variety of systems via techniques that include calorimetry (5, 6), oxygen consumption (30–33), and fluorescent metabolite measurements (30, 34, 35). Some of these studies capture the total energy expenditures of some biological system of interest and reveal intriguing mismatches between the measured dissipation and contribution such as the mechanical power. Accompanying theoretical efforts have reached similar conclusions (36, 37). Models have also even started to probe how biochemical energy sources might vary in space (38). Yet scarce experimental measurements impede their calibrated application to cells (35, 36). Further, these methods do not always have the spatial resolution that can isolate individual mechanisms, their contributions to the total dissipation, nor related gradients on the cellular or sub-cellular scale.

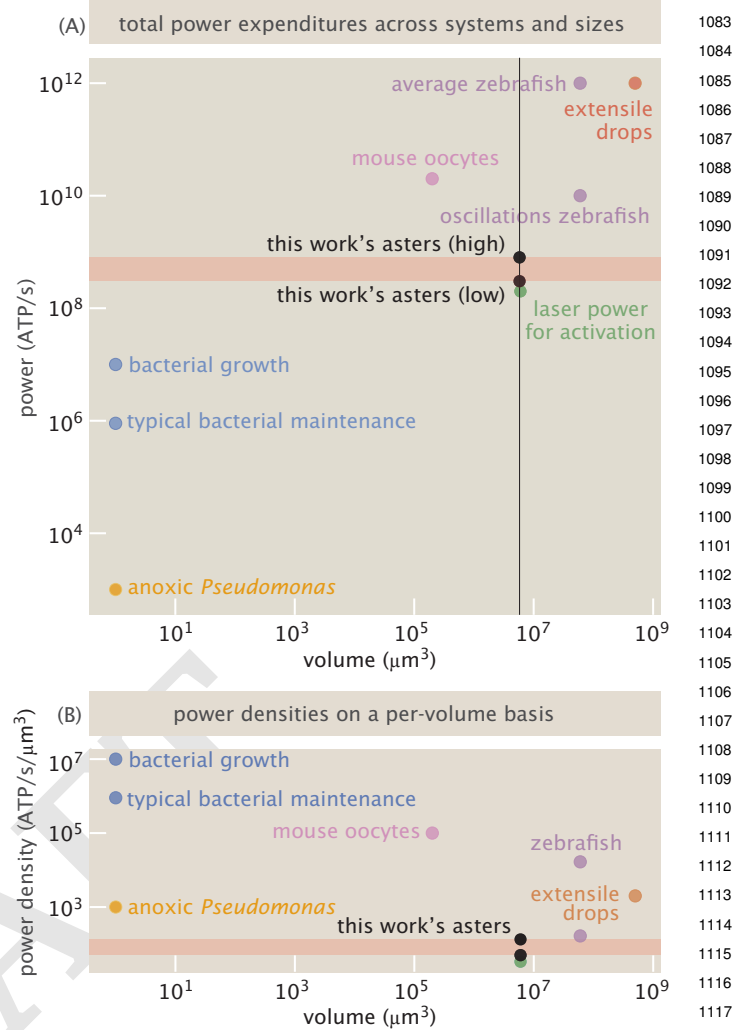
Having the ability to measure gradients is crucial since gradients are ubiquitous and functionally critical for processes such as body symmetry and division axes (29), division time, and cell size (26). These gradients resist thermodynamic equilibrium, requiring continuous energy input for their formation and maintenance. Motivated by the need to develop a physical understanding of how energy fluxes give rise to biological order, we developed

1019 an *in vitro* system of motor-microtubule assemblies to
 1020 ask how ATP consumption across space and time drives
 1021 the dynamics of structure formation. Within this frame-
 1022 work, ATP gradients emerge as a consequence of motor
 1023 protein activity. These gradients serve as a measurable
 1024 signature of the energetic mechanisms associated with the
 1025 rearrangement of microtubules during aster formation.

1026 In our *in vitro* system, such gradients emerge from
 1027 complex interactions between motors, microtubules, flows,
 1028 diffusion, and ATP/ADP competitive binding and are not
 1029 readily predicted from individual molecular properties. As
 1030 shown by our finite element modeling, this cooperation of
 1031 interactions can be quantitatively addressed by rigorous
 1032 modeling and real calibrated and thermodynamically-
 1033 meaningful measurements. These systems display gradi-
 1034 ents of order $\approx \text{few} \times 10^{-1} \mu\text{M}/\mu\text{m}$ over tens of microns—
 1035 hundreds of times steeper than celebrated developmental
 1036 gradients discussed in Figure 4.

1037 The differences between the magnitudes of the gradients
 1038 we measured and those observed in specific cellular and
 1039 organismal contexts (Figure 4), possibly come from a host
 1040 of additional complexities in those systems, which we
 1041 need to account for in future studies with living systems.
 1042 For example, in whole cells, there are metabolic energy
 1043 sources (including mitochondria) and sinks of energy other
 1044 than motors (filaments, ATP driven enzymes). It is a
 1045 fascinating, abiding question for our field, to ask whether
 1046 or how the conspiracy of both sources and sinks manifests
 1047 appreciable gradients (or not) of biochemical sources of
 1048 energy, as started in a beautiful recent work (38). We
 1049 can begin addressing these complexities by finite element
 1050 modeling, built upon our initial steps described here.
 1051 In terms of future studies, our analysis cautions that
 1052 the dissipation rate estimated by mean-field quantities
 1053 (such as average, not local, concentrations) can generically
 1054 overestimate the true macroscopic dissipation rate, as
 1055 we dissect further in the Supplementary Information
 1056 section S12 using an application of Jensen’s inequality.
 1057 This highlights the need for spatially resolved measure-
 1058 ments. To understand the nonequilibrium nature of whole
 1059 cells and their assemblies into tissues and organisms, we
 1060 acknowledge that an approach with at least a similar level
 1061 of rigor and scope as the one we have demonstrated needs
 1062 to be undertaken in these further contexts.

1063 Our experiments, calibrated in real molecular units,
 1064 resolve substantial dissipation rates, up to a few $\times 10^8$
 1065 ATPs per second ($> 10^{-17} \text{ W}/\mu\text{m}^3$), which for com-
 1066 parison, exceeds the average optical power densities
 1067 of our activation light of $10^{-17} \text{ W}/\mu\text{m}^3$. As illus-
 1068 trated in Figure 8, these power values measured in
 1069 contracting asters enjoy informative comparisons with
 1070 typical metabolic power expenditures of bacteria ($\approx 10^3$
 1071 ATP/s/cell in slow-metabolizing anaerobic *P. aeruginosa*
 1072 (39); $\approx 10^6$ ATP/s/cell (40, 41) in basal metabolism
 1073 to $\approx 10^7$ ATP/s/cell in exponential growth (42, 43)
 1074 in organisms like *E. coli*); mouse oocytes (approximately 10^5
 1075 ATP/ $\mu\text{m}^3/\text{s}$ totaling 2×10^{10} ATP/s/oocyte as inferred
 1076 from NADH turnover) (30); *Xenopus* embryos (on average,
 1077 approximately $\approx 1.5 \times 10^4$ ATP/ $\mu\text{m}^3/\text{s}$ in the two-cell-
 1078 stage to $\approx 2.1 \times 10^4$ ATP/ $\mu\text{m}^3/\text{s}$ at the cleavage stage as
 1079 inferred by calorimetry, with cell cycle oscillations about
 1080 one- hundred-fold smaller (6)); or other, extensile, active
 1081 matter microtubule systems (few $\times 10^3$ ATP/s/ μm^3) (5);
 1082



1109 **Fig. 8.** Comparisons of this work’s measured thermodynamic dissipations with those
 1120 of relevant organismal and active matter contexts. (A) Absolute power versus total
 1121 volume for embryos, microbial cells, and distinct experimental contexts. (B) The power
 1122 expenditures of these contexts measured on a per-volume basis. See SI §11 for
 1123 further details.

1124 see Supplementary Information section S11 for further
 1125 numerical discussion. Our experiments further report
 1126 ATP gradients develop in space (with steepnesses of order
 1127 at least a few tenths of micromolar per micron, sustained
 1128 over tens of microns).

1129 Beyond the aim of developing an energetic consump-
 1130 tion census, the formation of ATP gradients motivates
 1131 questions that ask: might these energetic gradients hold
 1132 significance to the cell? As we have highlighted in the
 1133 results, the energetic cost of maintaining a gradient is
 1134 comparable to energy expenditure we observe. Widely
 1135 celebrated cellular gradients, including the morphogen
 1136 gradient in *Drosophila* (29), which assigns the embryo’s
 1137 anterior-posterior body axis, the protein kinase gradient
 1138 in fission yeast (26), which regulates the timing of cellular
 1139 division, ras-related nuclear protein gradients (44), which
 1140 control spindle assembly, and others, are so beloved due to
 1141 how they drive cellular function. Perhaps, ATP gradients
 1142 have similar implications, localizing energy availability
 1143 to incite cellular activity in specified areas. Only with
 1144 measurements and mathematical models yielding space-
 1145

time understandings of energetic dissipation can one truly know how cells exists away from equilibrium.

Materials and Methods

Experimental, computational, and theoretical methods and results are described extensively in the supplementary information.

Data, Materials, and Software Availability. Data and analysis code used to generate this study's analyses are available open source; see <https://github.com/RPGroup-PBoC/am.atp>.

ACKNOWLEDGMENTS. We gratefully thank Frank Jülicher, Justin Bois, Xingbo Yang, Erwin Frey, Sara Mahdavi, Avi Flamholz, Peter Foster,

Catherine Ji, Victor Gomez, Vahe Galstyan, Mohammadamin Tajik, and members of the Phillips and Thomson groups for insightful discussions. We would also like to thank the David Van Valen and Rebecca Voorhees labs for providing resources for performing protein expression and purification. We thank the NIH for support through award numbers DP1OD000217 (Director's Pioneer Award) and NIH MIRA 1R35 GM118043-01 to RP. MT thanks The Moore-Simons Project on the Origin of the Eukaryotic Cell. AID and GLS were supported by the NSF Graduate Research Fellowship DGE-1745301. GLS thanks the NSF-Simons National Institute for Theory and Mathematics in Biology (NITMB) Fellowship supported via grants from the NSF (DMS-2235451) and Simons Foundation (MP-TMPS-00005320).

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